# 1 A Twofold Bernoulli Experiment with Complementary Outcomes in $\mathbb C$

Consider a Bernoulli experiment in which the outcome X can take two values:

X = 0 with complex probability a and X = 1 with complex probability b = 1 - a.

We perform this experiment twice independently. Thus, the probabilities for the first experiment are

$$Q(X_1 = 0) = a, \quad Q(X_1 = 1) = b,$$

and for the second experiment they are

$$Q(X_2 = 0) = a, \quad Q(X_2 = 1) = b.$$

**The Question.** Under what conditions on a and b does the twofold experiment yield, with probability one, complementary outcomes, i.e.

$$Q(X_2 = 1 - X_1) = 1?$$

In other words, we require that the probability of obtaining equal outcomes (denoted by p) and the probability of obtaining complementary outcomes (denoted by q) satisfy

$$p+q=1$$
 with  $q=1$ .

**Tree Diagram Analysis.** For two independent trials, the probability that the two outcomes are the same is given by

$$p = Q(X_1 = 0) \cdot Q(X_2 = 0) + Q(X_1 = 1) \cdot Q(X_2 = 1) = a^2 + b^2,$$

while the probability that the outcomes differ is

$$q = Q(X_1 = 0) \cdot Q(X_2 = 1) + Q(X_1 = 1) \cdot Q(X_2 = 0) = 2ab.$$

Since a + b = 1, it follows that

$$a^2 + b^2 = 1 - 2ab = 1 - q.$$

Thus, the two probabilities are related by

$$p = 1 - q$$
 and  $q = 2ab$ .

Finding a and b via the Quadratic Formula. If we view a and b as the two roots of a quadratic polynomial, then the polynomial with roots a and b is

$$t^{2} - (a+b)t + ab = t^{2} - t + \frac{q}{2} = 0.$$

The quadratic formula then gives:

$$t = \frac{1 \pm \sqrt{1 - 2q}}{2}.$$

We now distinguish several cases:

- If 1 2q > 0 (i.e.  $q < \frac{1}{2}$ ), then  $\sqrt{1 2q} \in \mathbb{R}$  and both a and b are real numbers. In particular, if q = 0 we have a = 1 and b = 0, which yields  $X_1 = X_2 = 0$  with certainty.
- If 1 2q = 0 (i.e.  $q = \frac{1}{2}$ ), then  $a = b = \frac{1}{2}$ .
- If 1 2q < 0 (i.e.  $q > \frac{1}{2}$ ), then  $\sqrt{1 2q} = i\sqrt{2q 1}$  is purely imaginary. In this case,

$$a = \frac{1}{2} + \frac{i\sqrt{2q-1}}{2}, \quad b = \frac{1}{2} - \frac{i\sqrt{2q-1}}{2}.$$

The Special Case q = 1. To have  $Q(X_2 = 1 - X_1) = q = 1$ , we must have p = 0. From the above relation, this occurs when

$$a^2 + b^2 = 0.$$

In the complex case, this forces

$$a = \frac{1}{2} + \frac{i}{2}, \quad b = \frac{1}{2} - \frac{i}{2}$$

Note that in this situation a + b = 1 and

$$|a| = |b| = \sqrt{\left(\frac{1}{2}\right)^2 + \left(\frac{1}{2}\right)^2} = \frac{1}{\sqrt{2}}.$$

Moreover, the phase difference between a and b is 90°. Thus, if one were to perform the experiment with these complex probabilities, the outcome would be perfectly anti-correlated; that is, if the first experiment yields  $X_1$ , then the second experiment yields  $X_2 = 1 - X_1$  with probability 1.

**Interpretation.** In the classical real-valued setting (q = 0), one can choose a = 1 and b = 0 so that both experiments yield the same outcome  $(X_1 = X_2 = 0)$ . However, to achieve perfect anti-correlation  $(X_2 = 1 - X_1 \text{ with certainty})$ , no such real numbers a and b exist. Instead, we must extend our probability assignments into the complex numbers. In this case, the choices

$$a = \frac{1}{2} + \frac{i}{2}$$
 and  $b = \frac{1}{2} - \frac{i}{2}$ 

yield q = 1, which formally corresponds to the quantum mechanical scenario of the EPR experiment with two entangled particles.

**Summary:** A twofold Bernoulli experiment with complex probabilities a and b = 1 - a yields perfectly complementary outcomes (i.e.  $X_2 = 1 - X_1$  with probability 1) if and only if

$$a = \frac{1}{2} + \frac{i}{2}$$
 and  $b = \frac{1}{2} - \frac{i}{2}$ 

This result highlights the need to extend the usual real-valued probabilities into the complex domain when modeling phenomena such as quantum entanglement.

### 2 Markov Chain Model for an "Entangled" Coin

We consider two distinct cases based on the parameter  $q \in [0, 1]$ :

#### 1. Case $0 \le q \le \frac{1}{2}$ : Ordinary Bernoulli Experiment

In this range, the quantity

$$a = \frac{1 + \sqrt{1 - 2q}}{2}$$

is a real number between 0 and 1 (specifically,  $0 \le a \le \frac{1}{2}$  if  $q \le \frac{1}{2}$ ). Consequently, we have a standard Bernoulli trial with probability *a* for Heads and b = 1 - a for Tails. By the Law of Large Numbers, if one performs *N* independent tosses, then approximately a fraction *a* of the outcomes will be Heads and a fraction *b* will be Tails.

#### 2. Case $\frac{1}{2} < q \le 1$ : Self-Entangled Coin with Complex Probabilities

Once  $q > \frac{1}{2}$ , the values of

$$a = \frac{1 \pm i\sqrt{2q-1}}{2}, \quad b = 1 - a$$

become complex (since  $\sqrt{1-2q}$  is purely imaginary). In this scenario, the coin exhibits "self-entangled" behavior in the sense that, if Heads occurs on one toss, then with probability  $q > \frac{1}{2}$  the next outcome is forced to be Tails (and conversely, Tails flips to Heads). Such a perfect complement flip is not observed in a usual physical coin.

To model this complementary-flip behavior in a classical (non-quantum) way, one may employ a two-state Markov chain with states H (Heads) and T (Tails). The transition matrix is

$$P = \begin{pmatrix} 1-q & q \\ q & 1-q \end{pmatrix},$$

so that with probability q the coin changes state on the next toss  $(H \to T, T \to H)$ , and with probability 1 - q it remains in the same state.

A simple calculation shows that the stationary distribution of this Markov chain is

$$\pi = (\frac{1}{2}, \frac{1}{2}).$$

Hence, by the Law of Large Numbers for Markov chains, one still obtains on average half Heads and half Tails over many tosses. However, the coin is "self-entangled" in the sense that it systematically prefers to *switch* (when  $q > \frac{1}{2}$ ) more than a fair coin would.

Thus, while the long-run frequencies might look similar (e.g. 50 % Heads, 50 % Tails), the dynamics and the underlying probabilities can be very different when  $q > \frac{1}{2}$ , requiring complex-valued probability amplitudes to formally describe the perfect anti-correlation scenario (q = 1).

## 3 A Quantum Mechanical Experiment with Perfect Anti-Correlation

A classic example in quantum mechanics that exhibits the desired behavior is the Einstein-Podolsky-Rosen (EPR) experiment performed with two spin- $\frac{1}{2}$  particles in the singlet state. In this experiment, the two particles are prepared in the entangled state

$$|\psi^{-}\rangle = \frac{1}{\sqrt{2}} \Big( |\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \Big),$$

where  $|\uparrow\rangle$  and  $|\downarrow\rangle$  represent the two possible outcomes of a spin measurement along a chosen axis. We associate these outcomes with the values X = 0 (say,  $\uparrow$ ) and X = 1 (say,  $\downarrow$ ), respectively.

Assume that we perform the following two independent measurements on the two particles along the same axis:

- The first measurement yields  $X_1$ .
- The second measurement is performed on the other particle.

Due to the perfect anti-correlation of the singlet state, if the first measurement yields

$$X_1 = 0,$$

 $X_2 = 1,$ 

 $X_1 = 1,$ 

then the second measurement will yield

with 100% probability. Conversely, if

then necessarily

$$X_2 = 0.$$

Thus, when we repeat the experiment on a large ensemble of such pairs, we observe that the outcome of the second measurement is always the complement of the outcome of the first measurement:

 $X_2 = 1 - X_1$  (with 100% probability).

This behavior illustrates the entangled nature of the quantum state and the non-classical correlations that arise from it.